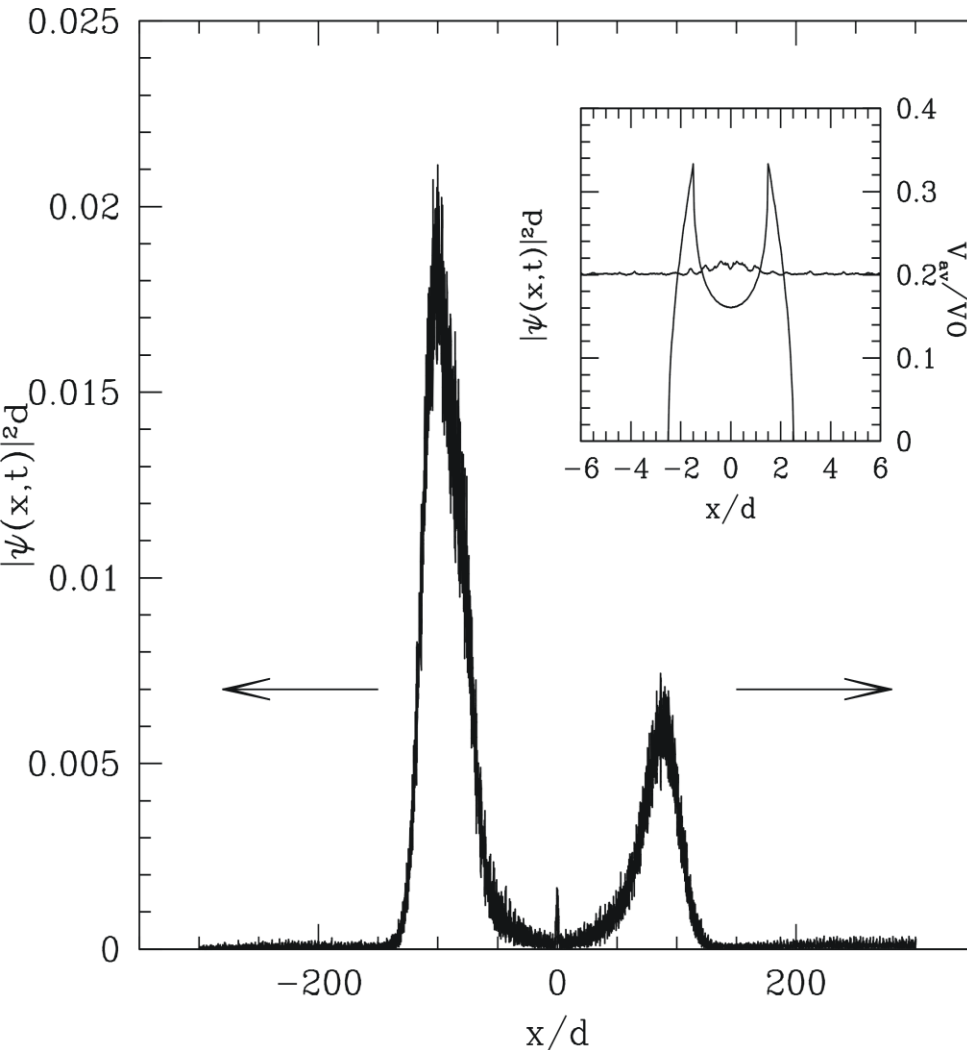


Analyze one effect at a time

- **Start with the case of noninteracting atoms at high frequency driving**
- **Proceed by tuning the frequency from very low to high**
- **Go back to high frequencies but switching on the nonlinear interactions**

The Physical Mechanism



❖ An atomic beam with $E_k = 0.2V_0$ near resonance and spread ΔE impinges on a light barrier of width d and height V_0 oscillating as

$$V(x,t) = V(x - l \cos(2\pi\nu t))$$

$$V(x) = V_0[\theta(x - d/2) - \theta(x + d/2)]$$

$$\nu = 10 \text{ KHz and } l = 2d$$

❖ A large portion of matter wave is transmitted well below barrier and a fraction still dwells inside the barrier

❖ Why (strictly true for $\nu \rightarrow \infty$): the time-averaged potential $V_{av} = 1/T \int V(x,t) dt$ allows for metastable states with energy E_0 and width Γ_0 possibly resonant with the incoming beam.

The Model

□ The problem can be considered as one-dimensional, since in the experiment the atomic beam travels along a waveguide. If the transverse confinement energy is much larger than all other energies, the wavefunction in the transverse direction is frozen

□ One-dimensional Gross-Pitaevskii equation, that is a nonlinear Schroedinger equation for the condensate wavefunction $\Psi(\mathbf{x}, t)$ along the direction of motion \mathbf{x}

$$i\hbar \frac{\partial \Psi(\mathbf{x}, t)}{\partial t} = \left[-\frac{\hbar^2 \nabla^2}{2m} + V(\mathbf{x}, t) + g |\Psi(\mathbf{x}, t)|^2 \right] \Psi(\mathbf{x}, t)$$

moving under the action of **the barrier $V(\mathbf{x}, t)$** oscillating at frequency ν (or the static **averaged potential $V_{av}(\mathbf{x})$**) \hbar

□ **Atomic interactions** enter the nonlinear term, with g modelled from the 3D value $g^{3D} = \frac{4\pi\hbar^2 a N}{m}$ to maintain the same level of average interactions

$$E_i = \frac{\int g |\Psi(x,t)|^4 dx}{\int |\Psi(x,t)|^2 dx} = \frac{\int g^{3D} |\Psi^{3D}(\mathbf{r},t)|^4 d\mathbf{r}}{\int |\Psi^{3D}(\mathbf{r},t)|^2 d\mathbf{r}}$$

For a gaussian wavepacket of oscillator length a_{\perp} and frozen in transverse direction the condition is satisfied by

$$g = \frac{g^{3D}}{\pi a_{\perp}^2}$$

Requirements and System Parameters

□ Requirements to observe a clean effect:

- ✓ $l > d$ to have double-barrier structure in V_{av}
- ✓ V_0 and l to have only one quasi-bound state in V_{av}
- ✓ $\Delta E = \Gamma_0$ (thus need to use a BEC)
- ✓ ν ? inverse tunnelling time to have high transparency

❑ System parameters are here tailored for sodium atoms from a transfer-matrix calculation and can be rescaled for other species

$$\begin{array}{cccc} d & V_0/h & E_0/h & \Gamma_0/h \\ 827nm & 4.2KHz & 0.90KHz & 0.06KHz \end{array}$$

❑ For the noninteracting atomic beam we have taken

$$\Delta E/h = 0.21KHz$$

❑ Reasonable set of parameters for the interacting atomic beam

$$\Delta E^*/h; 0.016 KHz = 0.0039 V_0 \quad g; 2g_0 \equiv 2 \frac{\hbar^2}{2md} \quad a_{\perp}; 1 \mu m \quad N=100$$

❑ Atomic velocities: 2 to 18 mm/s *i. e.* $0.1 < E_k < V_0 < 1$